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High fidelity and multiscale algorithms for collisional-radiative and nonequilibrium plasmas

*AFOSR Computational Math Review Meeting
July 2014*

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TOC



- **Problem description and motivation**
- **Translational equilibrium/non-equilibrium**
- **Level grouping**
- **Particle coalescence**
- **Multi-fluid equations**
- **Summary & future work**



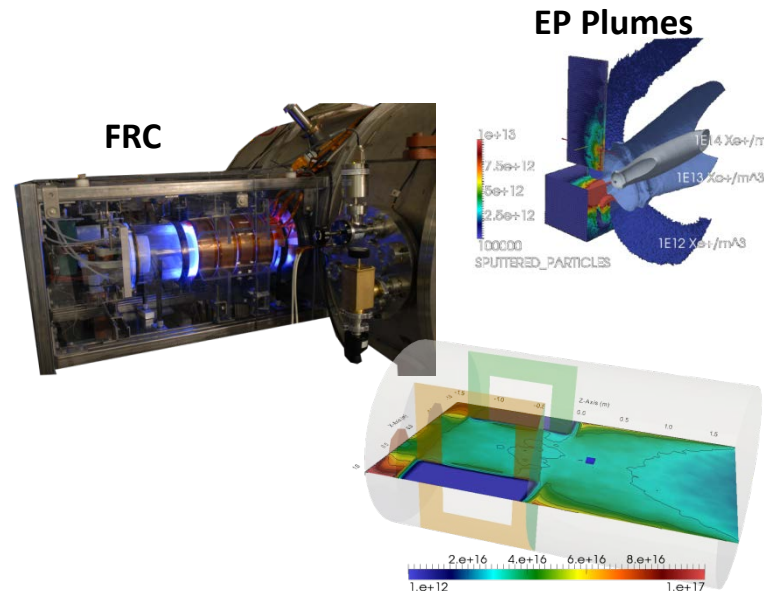
Spacecraft Plasma M&S



- Applications from Hall thrusters to plumes and gas discharges
- Complex physics: excitation/ionization, transport, radiative, material, etc.
- Multiple spatial-temporal and density scales.

Current focus:

Develop advanced multiscale algorithms for plasma M&S with highly non-equilibrium condition and collisional-radiative kinetics



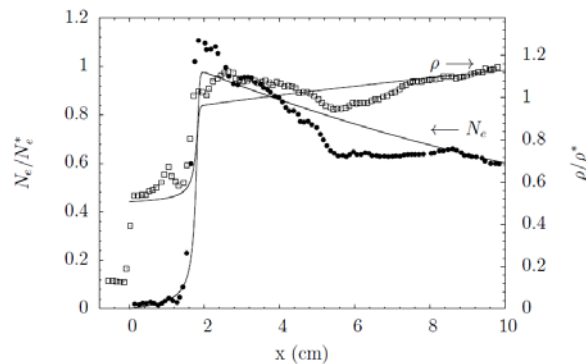
Chamber Environment



Collisional-Radiative kinetics



- **Non-equilibrium modeling of the atomic state distribution function (ASDF)**
 - Detailed state-to-state model of atomic transition, i.e., excitation and ionization
 - Rates derived based on ab initio cross section.
- **Examples: shocks**
- **Complications:**
 - Numerical **stiffness**



Atomic CR
Ar shock (31 levels)



Multiscale problem

- **Generic Multiscale problem:** $\frac{dy}{dx} = f_0(y, x) + \frac{1}{\varepsilon} f_1(y, x) \quad \varepsilon \rightarrow 0$
- **Various approaches possible:** DNS, multi-grid/AMR, HMM, equation-free, etc.

- Some similarities between approaches

- **Traditional approach: scale separation** $y = y_s + y_f$

- when $\varepsilon \rightarrow 0$, solve $\frac{dy_f}{dx} = \frac{1}{\varepsilon} f_1(y_s, y_f, x)$

- Use "relaxed" solution for $\frac{dy_s}{dx} = f_0(y_s, y_f^*, x)$

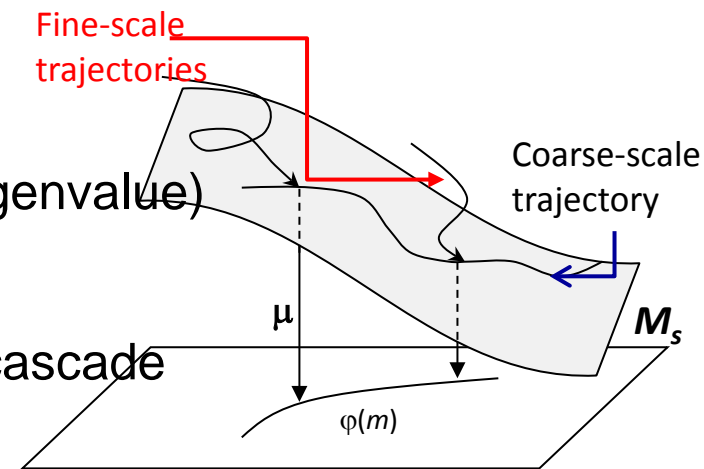
- Works slow manifold is attractive (neg. eigenvalue)

- **Complications:**

- Pos. eigenvalues. Ex: instability, inverse cascade

- Chaotic, stochastic fine scales, ...

- Non-separation of scales

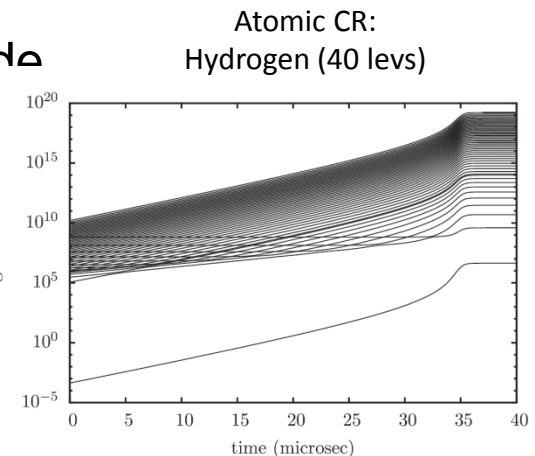




Multiscale problem

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 - Works slow manifold is attractive (neg. eigenvalue)
- **Complications:**
 - Pos. eigenvalues. Ex: instability, inverse cascade
 - Chaotic, stochastic fine scales, ...
 - Non-separation of scales: **CR kinetics**

15 orders of magnitude





Kinetic equation

- **Boltzmann equation:**

- $\varepsilon \ll 1$: fluid regime

- $\varepsilon \sim O(1)$: kinetic regime, need to resolve vfd

$$\partial_t f + \vec{v} \cdot \nabla_{\vec{x}} f = \frac{1}{\varepsilon} Q(f, f) + \dots$$

$$Q(f, f) = \int_{R^3} \int_{S^2} \sigma(|v - v_*|, \omega) [f(v')f(v_*') - f(v)f(v_*)] d\omega dv'$$

$$Q(f, f) = 0 \rightarrow f(v) = \frac{\rho}{(2\pi T)^{3/2}} \exp\left(-\frac{|v - u|^2}{2T}\right)$$

- For collisional plasma (fully and partially ionized): need to include Coulomb collisions (FP), excitation/ionization, CX, etc.

- **Need efficient algorithms for collisional-radiative processes in both regimes and **transitional** regime.**

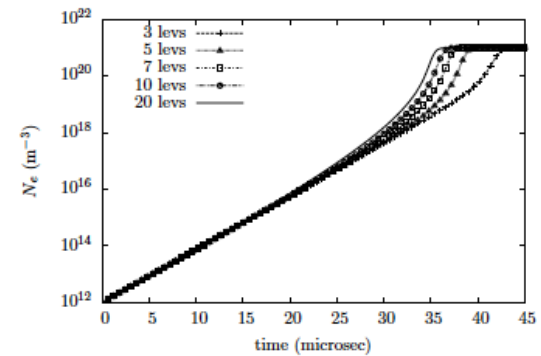


Translational Equilibrium

- **Moment method: derived by taking moment of the kinetic equations with fluid closure**
 - 5-moment model yields Euler/NS systems: multi-species, multi-temperature CR models.
 - Magnetized plasmas are often modeled with MHD with a hierarchy of descriptions: Ideal, resistive, Hall MHD.
 - Generalized model: multi-fluid (Braginskii, 1965)
- **CR kinetics = rate equations for each excited states**
 - Non-Boltzmann population of ASDF
 - Challenges: many states

$$\begin{aligned} \frac{dN_n}{dt} = & -\sum_{m>n} \alpha_{(m|n)} N_e N_n + \sum_{m>n} \beta_{(n|m)} N_e N_m + \sum_{m>n} A_{(n|m)} N_m \\ & + \sum_{m<n} \alpha_{(n|m)} N_e N_m - \sum_{m<n} \beta_{(m|n)} N_e N_n - \sum_{m<n} A_{(m|n)} N_n \\ & - \alpha_{(+|n)} N_e N_n + \beta_{(n|+)} N_+ N_e^2. \end{aligned} \quad (8)$$

$$\frac{dN_+}{dt} = \sum_n \alpha_{(+|n)} N_e N_n - \sum_n \beta_{(n|+)} N_+ N_e^2.$$



Atomic hydrogen (20 levels)

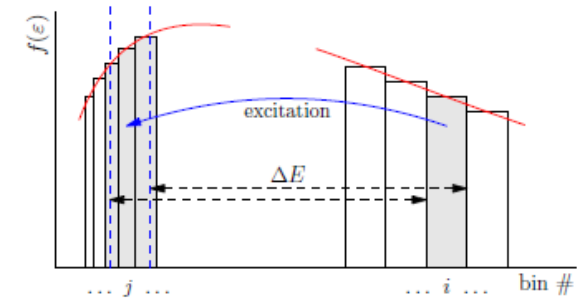


Translational Non-equilibrium



- **Not-too-far from equilibrium**

- Discretized EEDF yields rate equations for discrete elements (“bin”)
- DB enforced at microscopic level
- High-order, implicit and energy conserving
- More efficient compared to MCC.



$$\bar{n}_i = N_e \int_{\varepsilon_i}^{\varepsilon_i + \Delta\varepsilon} d\varepsilon f_e(\varepsilon)$$

- **Far from equilibrium**

- Particle methods with DSMC collisions
- Can resolve anisotropic vdf
- Drawback: slow convergence, reaction branching, singular rates, particles growth

$$\frac{d\bar{n}_i}{dt} = -N_l \bar{n}_i \sum_j \bar{k}_{(j|i)}^{\text{exc}} + N_u \sum_j \bar{n}_j \bar{k}_{(i|j)}^{\text{dex}}$$
$$\frac{d\bar{n}_j}{dt} = +N_l \sum_i \bar{n}_i \bar{k}_{(j|i)}^{\text{exc}} - N_u \bar{n}_j \sum_i \bar{k}_{(i|j)}^{\text{dex}}$$



Dynamical Regime

- **Examples:** thruster plumes, gas discharges, LPI, etc.
- **Require hybridization:** fluid/kinetic
- **Key challenges:** smooth, efficient and consistent transition
- **What we need:**
 - Multiscale statistics
 - Coarse-graining procedure for atomic state
 - Adaptive fluid-kinetic model
- **Current work:**
 - CR level grouping
 - Particle merging schemes
 - Multifluid model



Level grouping

- **CR modeling: level-grouping** $\rightarrow \mathcal{N}_n = N_{n0} \sum_{i \in n} \frac{N_i}{N_{n0}} \simeq \frac{N_{n0}}{g_{n0}} \underbrace{\sum_{i \in n} g_i e^{-\Delta E_i / T_n}}_{Z_n}$
 - Group effective rates of change
$$\frac{d\mathcal{N}_n}{dt} = -N_e \mathcal{N}_n \left[\sum_{m>n} \sum_{i \in n} \frac{g_i e^{-\Delta E_i / T_n}}{Z_n} \sum_{j \in m} \alpha_{(j|i)} + \sum_{m<n} \sum_{i \in n} \frac{g_i e^{-\Delta E_i / T_n}}{Z_n} \sum_{j \in m} \beta_{(j|i)} \right]$$
 - Internal structure of group is assumed Boltzmann (T_n)
 - Piecewise exponential
 - This does NOT mean the entire ASDF is Boltzmann!!
 - Group temperature must be determined \rightarrow additional conservation equation, e.g.:
$$\frac{d\mathcal{E}_n}{dt} = -N_e \mathcal{N}_n \left[\sum_{m>n} \sum_{i \in n} \frac{g_i e^{-\Delta E_i / T_n}}{Z_n} \sum_{j \in m} E_i \alpha_{(j|i)} + \sum_{m<n} \sum_{i \in n} \frac{g_i e^{-\Delta E_i / T_n}}{Z_n} \sum_{j \in m} E_i \beta_{(j|i)} \right]$$
 - Procedure?

– Solve:

$$\langle \Delta E \rangle_n (T_n^*) + C_v(T_n^*) \delta T_n^* = \frac{\Delta \mathcal{E}_n^{(k)}}{\mathcal{N}_n^{(k)}} \quad \text{with:}$$

$\xrightarrow{\text{iterated}}$

$\xrightarrow{\text{computed}}$

$\xrightarrow{\text{tabulated}}$

$\Delta \mathcal{E}_n = \sum_{i \in n} (E_i - E_{n0}) N_i = \mathcal{N}_n \langle \Delta E \rangle_n$
 $C_v(T_n) = \frac{d}{dT_n} \langle \Delta E \rangle_n$

– **However...** $\langle \Delta E \rangle_n \simeq o(\epsilon)$ where $\epsilon = e^{-\Delta E_1 / T_n} \rightarrow \delta T_n^* = o(\epsilon) / o(\epsilon)$
 $C_v(T_n) \simeq o(\epsilon)$



Level grouping

- **CR modeling: level-grouping**

- Other approaches?

- Sub-partitioning: lowest level n_0 and total \mathcal{N}_n (no need for \mathcal{E}_n)

$$\delta T_n^* \simeq \frac{T_n^{*2}}{\mathcal{Z}_n(T_n^*) \langle \Delta E \rangle_n(T_n^*)} \left[\frac{\mathcal{N}_n^{(k)}}{N_{n_0}^{(k)}} g_{n_0} - \mathcal{Z}_n(T_n^*) \right] = o(\epsilon)/o(\epsilon) \text{ ...fails}$$

- Sub-partitioning: lowest level n_0 and upper distribution \mathcal{N}_n ,

$$\delta T_n^* \simeq \frac{T_n^{*2}}{\mathcal{Z}'_n(T_n^*) \langle \Delta E \rangle_n(T_n^*)} \left[\frac{\mathcal{N}_n^{(k)}}{N_{n_0}^{(k)}} g_{n_0} - \mathcal{Z}'_n(T_n^*) \right] = o(\epsilon)/o(\epsilon) \text{ ...fails}$$

- Approximate \mathcal{Z}_n by expanding around mean energy: $\overline{\Delta E}_n = \frac{1}{g_n} \sum_{i \in n} g_i \Delta E_i$.

$$\mathcal{Z}_n(T_n) = e^{-\overline{\Delta E}_n/T_n} \sum_{i \in n} g_i \left[1 - \cancel{\frac{\delta_i}{T_n}} + \frac{1}{2} \frac{\delta_i^2}{T_n^2} + \dots \right] \quad \text{where } \delta_i \equiv \Delta E_i - \overline{\Delta E}_n$$

- With n_0 , \mathcal{N}_n partitioning: $1/\ln(1+\epsilon)$...fails

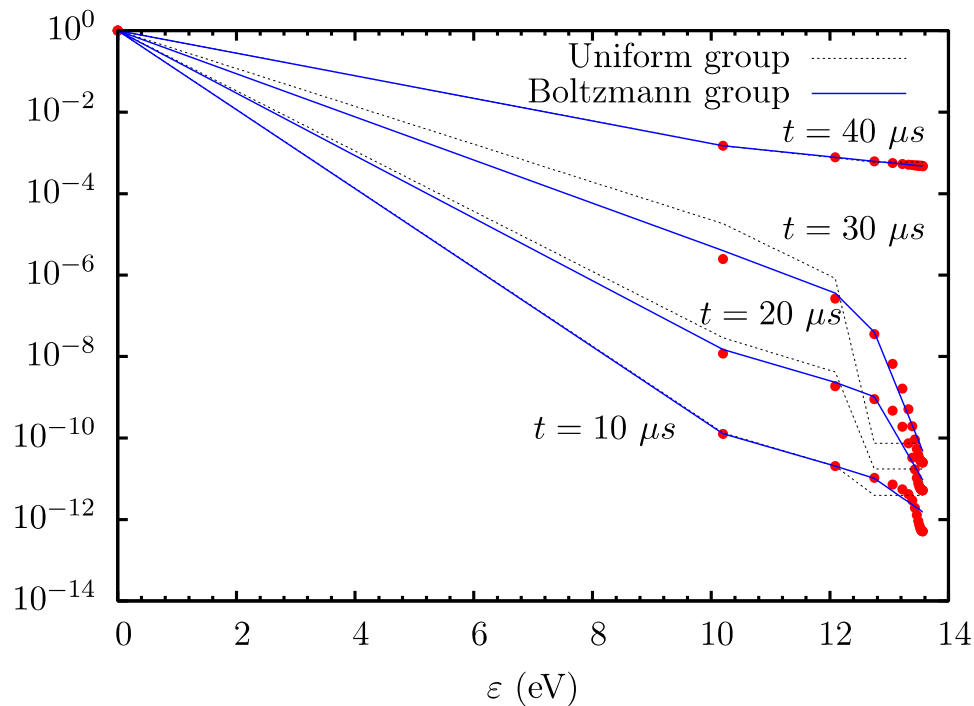
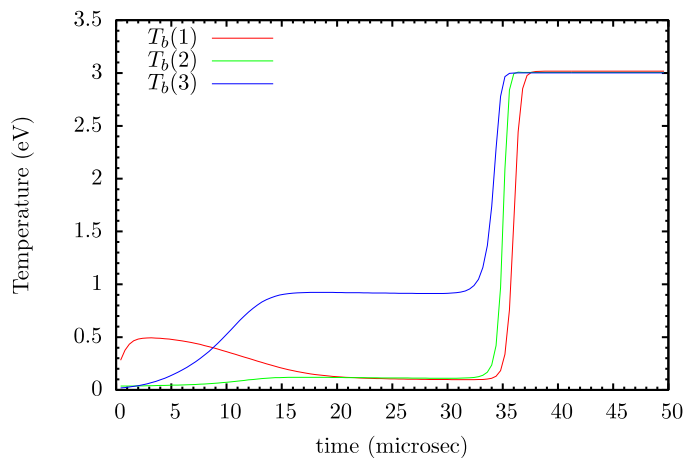
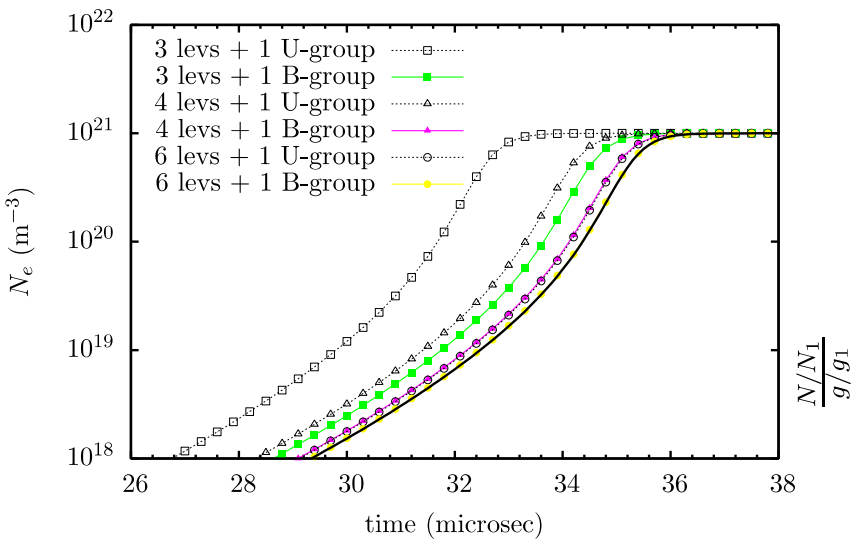
- With n_0 , \mathcal{N}_n partitioning: $1/\ln(\epsilon)$...succeeds!

- Improve with successive iterations... $\mathcal{Z}'_n(T_n) = \bar{g}'_n(T_n) e^{-\overline{\Delta E}'_n/T_n}$



Level grouping: 0D test

• Isothermal heat bath



Works very well (also tested in cooling regime).
Much better than uniform (standard) grouping.



Level grouping: Energy conservation

- CR modeling: level-grouping

- Look at uniform bins – 2 formulations:

~~$$\bar{\alpha}_{(m|n)}^E = \sum_{j \in m} \sum_{i \in n} \frac{g_i}{g_n} (E_j - E_i) \alpha_{(j|i)}$$~~

~~$$\bar{\beta}_{(m|n)}^E = \sum_{j \in m} \sum_{i \in n} \frac{g_j}{g_m} (E_j - E_i) \beta_{(i|j)}$$~~

$$\bar{\alpha}_{(m|n)}^E = (\bar{E}_m - \bar{E}_n) \sum_{j \in m} \sum_{i \in n} \frac{g_i}{g_n} \alpha_{(j|i)}$$

Energy conserving!

$$\bar{\beta}_{(m|n)}^E = (\bar{E}_m - \bar{E}_n) \sum_{j \in m} \sum_{i \in n} \frac{g_j}{g_m} \beta_{(i|j)}$$

- Right-form is product of Energy x $d\mathcal{N}_n/dt$...

- Conservation follows from definition $\mathcal{E}_n = \bar{E}_n \mathcal{N}_n$

- Same principle can be applied to Boltzmann groups:

- Start with:

$$\frac{d\Delta\mathcal{E}_n}{dt} \equiv \sum_{i \in n} \Delta E_i \frac{dN_i}{dt} = \frac{d}{dt} (\mathcal{N}'_n \langle \Delta E \rangle_{n'}) = \langle \Delta E \rangle_{n'} \frac{d\mathcal{N}'_n}{dt} + \mathcal{N}'_n \frac{d\langle \Delta E \rangle_{n'}}{dt}$$

$\Delta(\text{overall group population}) : \mathcal{N}_n$
 $\Delta(\text{internal structure}) := C_{v,n'} \frac{dT_n}{dt}$

- Express in terms of conserved variables:

$$\frac{d\mathcal{E}_n}{dt} = [E_{n_0} - \omega_{n'}] \frac{dN_{n_0}}{dt} + [E_{n_0} + \langle \Delta E \rangle_{n'} + \xi_{n'}] \frac{d\mathcal{N}'_n}{dt} \quad \text{with} \quad \xi_{n'} = \frac{C_{v,n'} T_n^2}{(\Delta E'_n + T_n^2 \frac{d \ln \bar{g}_n}{dT_n})} \quad \text{and} \quad \omega_{n'} = \xi_{n'} \frac{\mathcal{N}_{n'}}{N_{n_0}}$$



Level grouping: Energy conservation

- CR modeling: level-grouping

- Finally...Procedure shown to be equivalent to replacing energies by “effective” (condition-dependent) values (\approx EOS)

Excitation:

$$\tilde{\alpha}_{(m_0|n_0)}^E = (\tilde{E}_{m_0} - \tilde{E}_{n_0}) \cdot \tilde{\alpha}_{(m_0|n_0)}$$

$$\tilde{\alpha}_{(m'|n_0)}^E = (\tilde{E}_{m'} - \tilde{E}_{n_0}) \cdot \tilde{\alpha}_{(m'|n_0)}$$

$$\tilde{\alpha}_{(m_0|n')}^E = (\tilde{E}_{m_0} - \tilde{E}_{n'}) \cdot \tilde{\alpha}_{(m_0|n')}$$

$$\tilde{\alpha}_{(m'|n')}^E = (\tilde{E}_{m'} - \tilde{E}_{n'}) \cdot \tilde{\alpha}_{(m'|n')}$$

Ionization:

$$\tilde{\alpha}_{(+|n_0)}^E = (I_H - \tilde{E}_{n_0}) \cdot \tilde{\alpha}_{(+|n_0)}$$

$$\tilde{\alpha}_{(+|n')}^E = (I_H - \tilde{E}_{n'}) \cdot \tilde{\alpha}_{(+|n')}$$

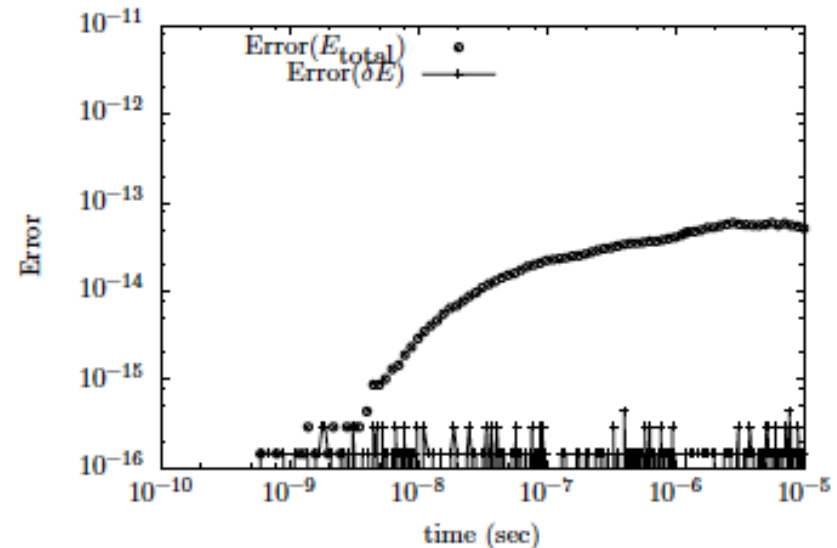
with

$$\tilde{E}_{n_0} = E_{n_0} - \omega_{n'}$$

$$\tilde{E}_{n'} = E_{n_0} + \langle \Delta E \rangle_{n'} + \xi_{n'}$$

- NOW**, energy is conserved

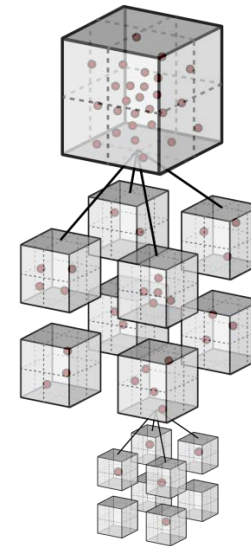
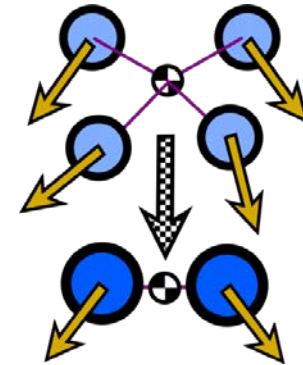
(down to round-off)





Particle Merging Schemes

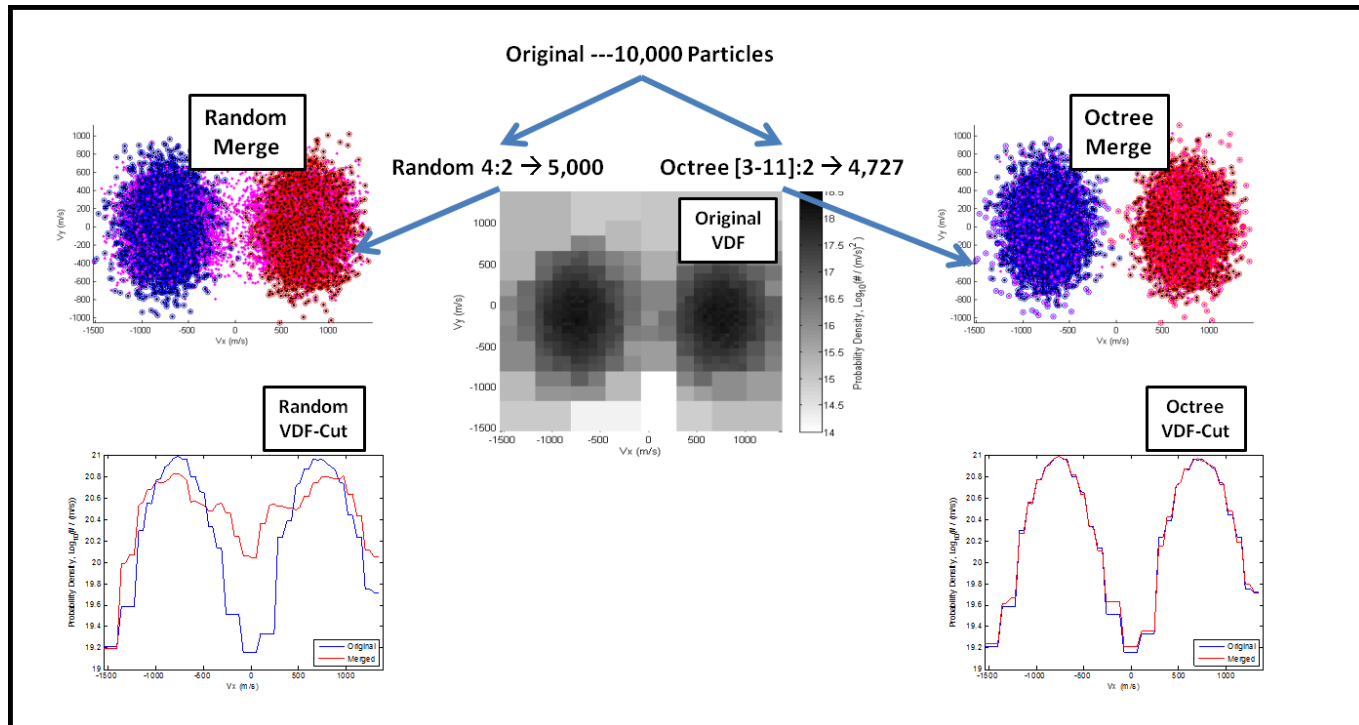
- The schemes consists of merging 3 or more particles to 2. Mass, momentum and kinetic energy are exactly conserved; Electrostatic energy also conserved in physical space.
- Split analogously defined by merging only fractions of original particles.
- To inhibit thermalization, an octree in velocity space is used so that only near neighbor particles are merged.
- Higher-moment conserving schemes obtained with increased number of merge result particles generated.





Examples

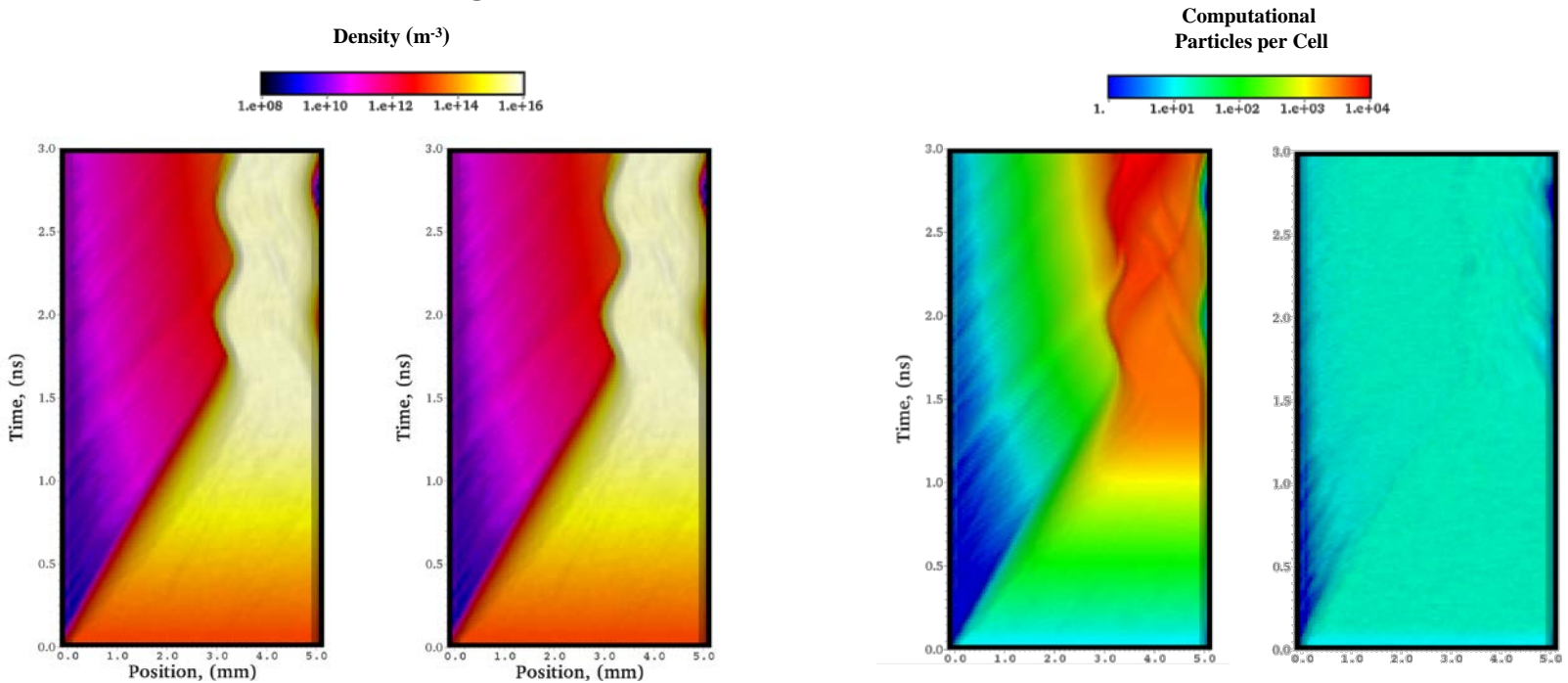
- from 0D ...





Examples

- from 0D to 1D3V...
 - Gas discharge with ionization.

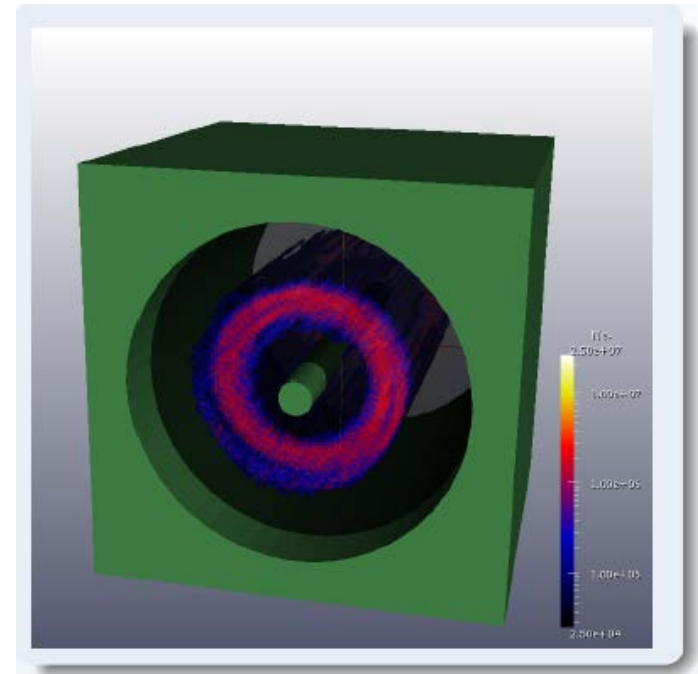




Examples



- from 0D to 1D3V and 3D3V ...
 - Annular test case
 - More test cases underway





Multi-fluid equations

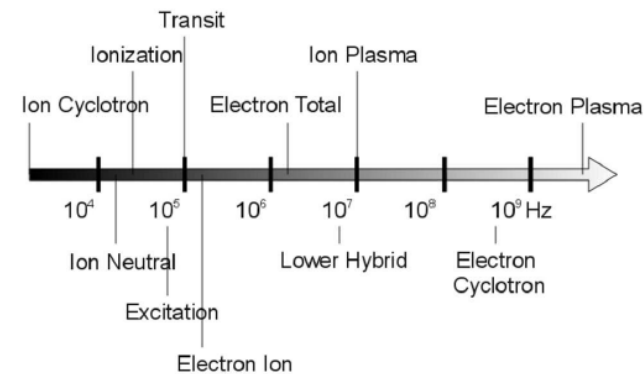
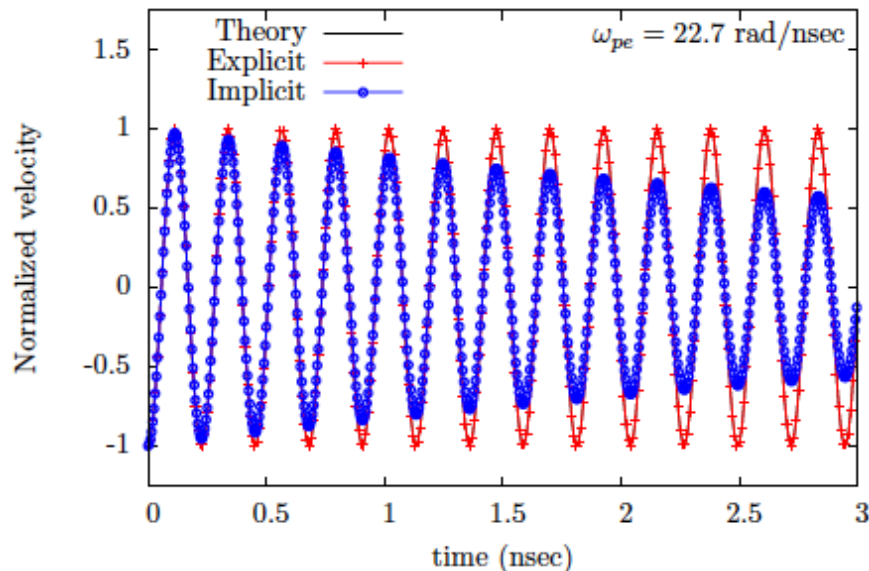
- **Multi-fluid model:**

- Develop series of models for variable conditions:

Electrostatic → + Magnetostatic → Electro-magnetic → 3-fluid

- Use implicit models to eliminate constraint of sequence of fast time scales: c , v_e , $\omega_{ps} = \sqrt{\frac{n_s q_s^2}{\epsilon_0 m_s}}$ $\omega_{cs} = \frac{q_s B}{m_s}$

- Price to pay: lack of resolution misses physics...





Multi-fluid equations

- **Multi-fluid model:**

- 5-moment:

$$\partial_t \rho_s + \nabla \cdot (\rho_s \mathbf{u}_s) = \omega_s^\rho$$

$$\partial_t (\rho_s \mathbf{u}_s) + \nabla \cdot (\rho_s \mathbf{u}_s \mathbf{u}_s + \mathbb{P}_s) = Z_s e n_s (\mathbf{E} + \mathbf{u}_s \times \mathbf{B}) + \mathbf{f}_s^m$$

$$\partial_t \epsilon_s + \nabla \cdot (\mathbf{u}_s \epsilon_s + \mathbb{P}_s \cdot \mathbf{u}_s) + \nabla \cdot \mathbf{q}_s = \mathbf{j}_s \cdot \mathbf{E} + \omega_s^\epsilon$$

- Add Maxwell's equations:

$$\begin{aligned} \nabla \times \mathbf{E} &= -\partial_t \mathbf{B} & \nabla \cdot \mathbf{E} &= \frac{e}{\epsilon_0} (Z_i n_i - n_e) \\ \nabla \times \mathbf{B} &= \mu_0 \mathbf{j} - \frac{1}{c^2} \partial_t \mathbf{E} & \nabla \cdot \mathbf{B} &= 0 \end{aligned}$$

- Add collisions:

- Elastic – Bragiinski terms

- Inelastic – warning! Rates depend on both T and relative velocity

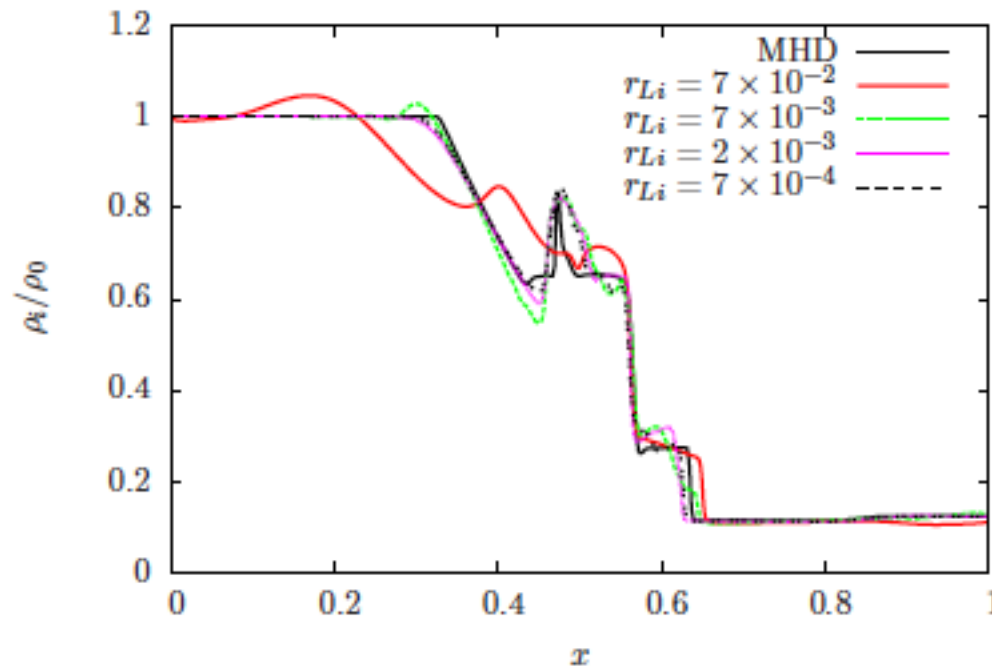
$$k_i = n_n n_e \int \int f_n f_e g \sigma_i''(g; \Omega_1, \Omega_2) d\Omega_1 d\Omega_2 d^3 v_n d^3 v_e \quad \longrightarrow \quad k_i = k_i(T_e, |\mathbf{u}_n - \mathbf{u}_e|)$$

- Multi-fluid CR model from fundamental principles being developed (incl. detailed balance)



Multifluid equations

- **Electromagnetic shock: generalized Brio-Wu¹**
 - FV with WENO reconstruction and RK3

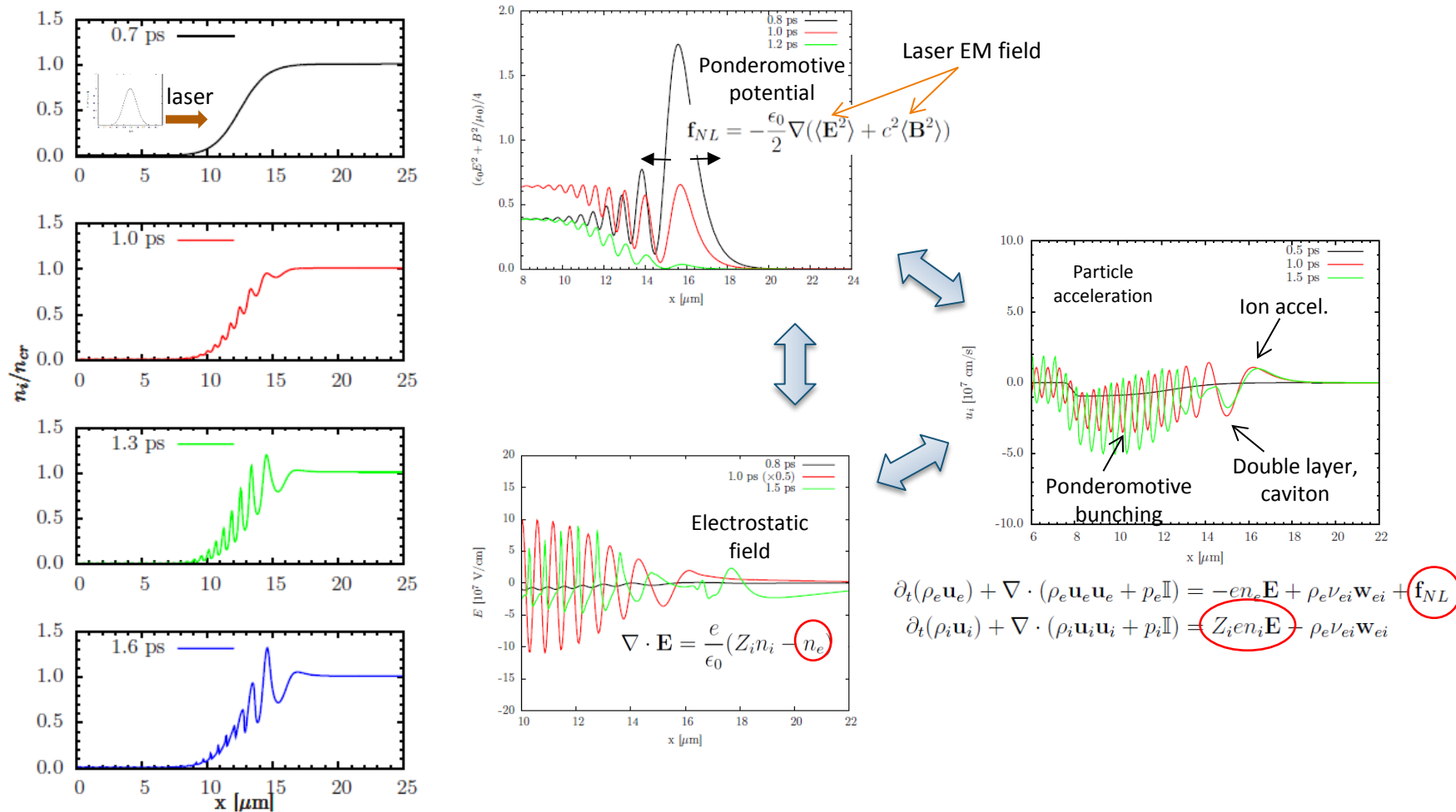


¹Shumlak & Loverich, JCP 2003



Laser-Plasma Interaction

- Ion acceleration due to ponderomotive force





Conclusion



- **Multiscale algorithms for nonequilibrium flows with CR kinetics**
 - Particle merge/split for particle management, efficient sampling, inelastic collisions ...
 - Level grouping schemes of electronic states, for dynamical coarse-graining of ASDF.
 - Multifluid equations to efficiently capture electron “hydrodynamics”
- **Ongoing works:**
 - High-order particle merging schemes
 - Multi-D simulation with level grouping
 - Modeling of inelastic collisions in multifluid